

Problem Set A2.

Kuba Skwarek

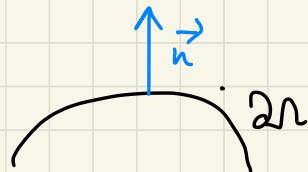
Problematik.

Divergence theorem : $F: \mathbb{R}^n \rightarrow \mathbb{R}^n$, $F \in C^1$,
 $\rightarrow F(x) = (F_1(x), \dots, F_n(x))$.

$$\begin{aligned}\rightarrow \operatorname{div} F(x) &:= \partial_{x_1} F_1(x) + \dots + \partial_{x_n} F_n(x), = \\ &= \nabla \cdot F = (\partial_{x_1}, \dots, \partial_{x_n}) \cdot F.\end{aligned}$$

$$\rightarrow \int_S \operatorname{div} F(x) dx = \int_{\partial \Omega} \langle F, n \rangle ds$$

\mathcal{S} $\partial \Omega$ ||
 $F \cdot n$



$$(A) \quad \Delta u = \operatorname{div}(\nabla u)$$

$$(\sum \partial_{x_i}^2 u = \sum \partial_{x_i} (\partial_{x_i} u)).$$

(B) Using divergence theorem

$$\int_S \Delta u = \int_{\partial \Omega} \langle \nabla u, n \rangle = \int_{\partial \Omega} \frac{\partial u}{\partial n} \nabla u \cdot n$$

Problem A2

Using div thm with $F = v \nabla u$ we have

$$\begin{aligned} \operatorname{div}(v \nabla u) &= \sum_{i=1}^n \partial_{x_i} (v \partial_{x_i} u) = \sum_{i=1}^n \partial_{x_i} v \cdot \partial_{x_i} u + \\ &+ v \sum_{i=1}^n \partial_{x_i}^2 u = \nabla v \cdot \nabla u + v \cdot \Delta u. \quad \square. \end{aligned}$$

Problem A3

Use Problem 2 twice.

$\nearrow j\text{-th}$

Problem A4

We consider $F = (0, 0, \dots, 0, u v, 0, \dots, 0)$. Then, $\operatorname{div} F = D_j u \cdot v + u \cdot D_j v$. Moreover,

$$\langle F, n \rangle = u v \cdot n_j.$$

Problem B1

\uparrow

Problem B2

(as martingales!!!)

Let u be harmonic. Then

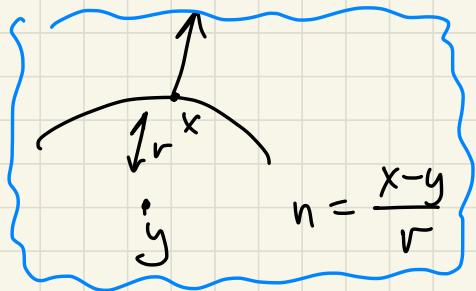
$$\Delta \psi(u) = \sum_{i=1}^n \partial_{x_i}^2 \psi(u) = \sum_{i=1}^n \partial_{x_i} \left[\psi'(u) \partial_{x_i} u \right] =$$

$$= \sum_{i=1}^n \underbrace{\psi'(u) \partial_{x_i}^2 u}_{\geq 0} + \underbrace{\psi''(u) |\nabla u|^2}_{\geq 0} \geq 0. \quad \square.$$

Problem B3

u is subharmonic i.e. $\Delta u \geq 0$. For all $B_r \subset \mathbb{R}^n$

$$0 \leq \int_{B_r(y)} \Delta u = \int_{\partial B_r(y)} \frac{\partial u}{\partial n} dS = \int_{\partial B_r(y)} \nabla u \cdot n dS \Rightarrow$$



$$\int_{\partial B_r(y)} \nabla u(x) \cdot \frac{x-y}{r} dS(x) \geq 0.$$

Now, consider $\phi(r) := \int_{\partial B(y, r)} u(x) dS$

$$\begin{aligned} d(r) &= \frac{1}{\omega(n)} r^{n-1} \int_{\partial B(y, r)} u(x) dS(x) \\ &\quad \uparrow \\ &\quad \frac{x-y}{r} = z \quad x = zr + y \\ &\quad dz = r^{n-1} d_z \end{aligned}$$

$$= \frac{1}{\omega(n) r^{n-1}} \int_{\partial B(0, 1)} u(zr + y) r^{n-1} dS(z)$$

$$= \int_{\partial B(0, 1)} u(zr + y) dS(z).$$

$$\phi(v) = \int_{B(0,1)} u(zr+y) dS(z)$$

$$\phi'(v) = \int_{B(0,1)} Du(zr+y) \cdot z dS(z) =$$

$$= \frac{1}{\alpha(n)} \int_{B(0,1)} Du(zr+y) \cdot z dS(z) \quad \begin{aligned} x &= zr+y \\ dx &= r^{n-1} dz \end{aligned}$$

$$= \int_{B(y,r)} \nabla u(x) \cdot \frac{x-y}{r} dS(z) \geq 0$$

so ϕ is non-decreasing in v . But $\phi(v) \rightarrow u(y)$
as $r \rightarrow 0$. So $\int_{B(y,r)} u(x) dS \geq u(y)$.

$$\text{To see 2nd formula } \int_{B(y,r)} u(x) dx =$$

$$= \int_0^r \int_{\partial B(y,w)} u(x) dS(x) dw = \int_0^r \alpha(n) w^{n-1} \int_{\partial B(y,w)} u(x) dS dw$$

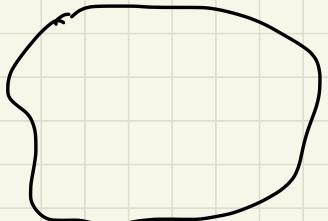
$$\geq u(y) \int_0^r \alpha(n) w^{n-1} dw = u(y) n \alpha(n) r^n.$$

□.

For superharmonic function $u(y) \geq \int_{\partial B_R(y)} u(x) d\mathcal{L}(x)$,
 For harmonic functions we have equality.

Problem B4

u is subharmonic in Ω and u attains its max inside Ω in some $x_0 \in \text{int } \Omega$. Then, for any $B_R(x_0)$



$$u|_{B_R(x_0)} = u(x_0) \text{ by B3.}$$

We cover Ω with such balls to conclude the proof.

Problem B5

Clearly, u attains its max. somewhere in $\overline{\Omega}$. If this happens inside $\Rightarrow u$ is constant in $\overline{\Omega}$, including the boundary. For harmonic function

$$\sup_{\Omega} u(x) = \sup_{\partial\Omega} u(x), \quad \inf_{\Omega} u(x) = \inf_{\partial\Omega} u(x).$$

Problem B6 $\Delta u \geq 0$ so u is subharmonic, i.e. u attains sup on the boundary $\Rightarrow u \leq 0$. If $u \geq 0$ then $u=0$ (and this is a solution).

Problem C1

If there were 2 solutions u_1, u_2 then

$$\begin{cases} \Delta(u_1 - u_2) = 0 & \text{in } \Omega, \\ u_1 - u_2 = 0 & \text{on } \partial\Omega. \end{cases}$$

Then $u_1 - u_2$ attains its max and min on $\partial\Omega$

so $u_1 - u_2 = 0$. \square .

Problem C2

Use A2 to $v = u = u_1 - u_2$

$$\int_{\Omega} (u_1 - u_2) \Delta (u_1 - u_2) + \int_{\Omega} |\nabla (u_1 - u_2)|^2 = \int_{\Omega} (u_1 - u_2) \frac{\partial u_1}{\partial n}$$

||
0
or ||
0 on $\partial\Omega$

$$\Rightarrow \int_{\Omega} |\nabla (u_1 - u_2)|^2 = 0 \Rightarrow u_1 - u_2 \text{ is constant.}$$

Use value at $\partial\Omega$ to conclude $u_1 = u_2$. \square .

Problem C3



Problem C4

If u_1, u_2 solve $\Delta u_1 = f, \Delta u_2 = f \Rightarrow$

$$\Delta(u_1 - u_2) = 0 \quad \text{in } \mathbb{R}^n$$

$\Rightarrow u_1 - u_2$ is harmonic in \mathbb{R}^n

$\Rightarrow u_1 - u_2 = C$ constant by Liouville Theorem

(so uniqueness is up to a constant C).

Probleme D1

$$\Phi(x) = \begin{cases} \frac{1}{n(2-n)\omega_n} |x|^{2-n} & n > 2 \\ \frac{1}{2\pi} \log|x| & n = 2 \end{cases}$$

$$|x| = \sqrt{x_1^2 + \dots + x_n^2}$$

$$D_i |x| = \frac{1}{2} \frac{1}{\sqrt{x_1^2 + \dots + x_n^2}} 2x_i = \frac{x_i}{|x|}$$

$$D_i |x|^{2-n} = (2-n) |x|^{1-n} \frac{x_i}{|x|} = (2-n) |x|^{-n} x_i$$

$$D_i \log|x| = \frac{1}{|x|} \frac{x_i}{|x|} = \frac{x_i}{|x|^2}$$

$$D_i \underline{\Phi}(x) = \begin{cases} \frac{1}{n\omega_n} \frac{x_i}{|x|^n} & n < 2 \\ \frac{1}{2\pi} \frac{x_i}{|x|^2} & n = 2 \end{cases} \leq \frac{1}{n\omega_n} |x|^{1-n}$$

$$D_{ij} \underline{\Phi}(x) = \frac{1}{n\omega_n} \frac{\delta_{ij} |x|^n - x_i n |x|^{n-1} \frac{x_j}{|x|}}{|x|^{2n}} =$$

$$= \frac{1}{n\omega_n} \frac{\delta_{ij} |x|^n - n x_i x_j |x|^{n-2}}{|x|^{2n}} \leq \begin{cases} i \neq j & \frac{1}{\omega_n |x|^n} \\ i = j & \frac{1}{\omega_n |x|^n} \end{cases}$$

Problem P2

Fundamental solution in 1D:

$u'' = 0$ and u is radially symmetric

u is linear, $u = -\frac{1}{2} |x|$.

$$u'(x) = \begin{cases} -\frac{1}{2} & x > 0 \\ +\frac{1}{2} & x < 0 \end{cases}$$

$"u" = -\delta_0$

meaning: jump at $x=0$.

Three meanings of $u''(x) = -\delta_0$

1) Function has jump of size 1 at one particular point.

$$2) \int f(x) d u''(x) = -f(0)$$

so this is measure $-\delta_0$.

3) This is linear functional on space of cont. functions:

Problem D3 "Δ $\underline{\Phi}(x) = -\delta_0(x)$ ".

$$\text{i.e. } \int f(y) \Delta \underline{\Phi}(y) dy = - \int f(x) \delta_0(x) dx \\ = -f(x).$$

Green's 2nd identity:

$$\int_{\Omega} (\nabla \Delta u - \Delta v u) = \int_{\partial\Omega} \left(\nabla \frac{\partial u}{\partial n} - \frac{\partial v}{\partial n} u \right)$$

We apply this with $v(y) = \underline{\Phi}(y-x)$, x fixed.
 $u(y) = u(y)$

$$\text{Then } \int_{\Omega} \Delta \underline{\Phi}(y-x) \cdot u(y) dy = \int_{\Omega} u(y) \delta_x dy \\ = 0 \Leftrightarrow y=x \\ = -u(x).$$

$$u(x) = - \int_{\Omega} \underline{\Phi}(y-x) \Delta u(y) dy + \int_{\partial\Omega} \underline{\Phi}(y-x) \frac{\partial u}{\partial n} dS(y) \\ - \int_{\Omega} \frac{\partial \underline{\Phi}}{\partial n}(y-x) \cdot u(y) dS(y).$$

The rigorous proof in the lecture.

Problem D4 We know:

$$u(x) = - \int_{\Omega} \Phi(y-x) \Delta u(y) dy + \int_{\partial\Omega} \Phi(y-x) \frac{\partial u}{\partial n} dS(y)$$
$$- \int_{\partial\Omega} \frac{\partial \Phi}{\partial n}(y-x) \cdot u(y) dS(y).$$

troublemaker

Fix x and consider corrector solving for fixed x :

$$\begin{cases} \Delta \varphi^x = 0 & \text{in } \Omega \\ \varphi^x = \Phi(x-y) & \text{on } \partial\Omega \end{cases}$$

$\Phi(y-x)$

Then

$$0 = - \int_{\Omega} \varphi^x(y) \Delta u(y) dy + \int_{\partial\Omega} \varphi^x(y) \frac{\partial u}{\partial n} dS(y)$$
$$- \int_{\partial\Omega} \frac{\partial \Phi}{\partial n}(y-x) u(y) dS(y)$$

Hence

$$u(x) = - \int_{\mathcal{L}} \left[\Phi(y-x) - \zeta^x(y) \right] \Delta u(y) dy$$

$$\sim \int_{\partial \mathcal{L}} \frac{\partial}{\partial n} \left[\Phi(y-x) - \zeta^x(y) \right] u(y) dy.$$

This motivates GREEN FUNCTION for \mathcal{L} :

$$G(x,y) = \Phi(y-x) - \zeta^x(y).$$

So far, we know that if $u \in C^2(\mathcal{L}) \cap C(\bar{\mathcal{L}})$ solves

$$\begin{cases} -\Delta u = f & \mathcal{L} \\ u = g & \partial \mathcal{L} \end{cases} \quad \text{then}$$

$$\Delta G(x,y) = 0 \quad x \neq y.$$

$$u(x) = \int_{\mathcal{L}} G(x,y) f(y) dy - \int_{\partial \mathcal{L}} \frac{\partial G(x,y)}{\partial n_y} g(y) dy.$$

BUT DOES IT EXIST?

Now, we recall some facts from the lecture:

① $G(x,y) = G(y,x)$.

② If $\mathcal{L} = B_r(0)$ then one can compute G to

obtain G s.t. $\frac{\partial G(x,y)}{\partial n_y} = -\frac{r^2 - |x|^2}{n \alpha_n r} \frac{1}{|x-y|^n}$.

In particular, when $\begin{cases} -\Delta u = 0 \\ u = g \end{cases}$ we know

that

$$u(x) = \int_{\partial B(0,r)} \frac{r^2 - |x|^2}{n \alpha_n r} \frac{g(y)}{|x-y|^n} dS(y).$$

Problem D5

Using Green function one can obtain formula for solution to $\begin{cases} \Delta u = 0 \\ u = g \end{cases}$ $\Omega = B_r(0)$.

$$u(x) = \frac{r^2 - |x|^2}{n \alpha_n r} \int_{\partial B(0,r)} \frac{g(y)}{|x-y|^n} dy$$

$$= - \int \frac{\partial G(x,y)}{\partial n_y} g(y) dy$$

$$\frac{\partial G(x,y)}{\partial n_y} = -\frac{r^2 - |x|^2}{n \alpha_n r} \frac{1}{|x-y|^n}.$$

We will prove:

$$(a) \quad u \in C^\infty(B_r(0))$$

$$(b) \quad \Delta u = 0 \quad \text{in } B_r(0)$$

$$(c) \quad \forall_{x_0 \in \partial B_r(0)} \lim_{\substack{x \rightarrow x_0 \\ x \in B_r(0)}} u(x) = g(x_0).$$

Proof of (a) and (b) :

For fixed x , the map $y \mapsto G(x,y)$ is harmonic except $y=x$. As $\overline{G(x,y)} = G(y,x)$, the map $x \mapsto G(x,y)$ is harmonic, except $x=y$. Hence $x \mapsto \frac{\partial G}{\partial y_x}(x,y)$ is harmonic. Hence

$$\Delta_x \frac{\partial G}{\partial y_x}(x,y) = 0 \quad x \in B_r(0), y \in \partial B_r(0)$$

Directly from the formula we see that

$\left\| \Delta_x \frac{\partial G}{\partial y_x} \right\|_{L^\infty(B_{r(1-\varepsilon)}(0))}$ is bounded

so DCT shows $\forall x \in B_{r(1-\varepsilon)}(0) \quad \Delta u(x) = 0$.
 Finally, we send $\varepsilon \rightarrow 0$ to conclude the proof.

Proof of (c):

Note that $u(x) = 1$ is harmonic so applying theory above

$$1 = u(x) = \int_{\partial B_r} -\frac{\partial G}{\partial n_y}(x, y) \cdot 1 \, dy = - \int_{\partial B_r} \frac{\partial G}{\partial n_y}(x, y) \, dy$$

Hence, fix $\varepsilon > 0$, choose δ $|g(x_0) - g(y)| \leq \varepsilon$ for $|x_0 - y| \leq \delta$
 If $|x_0 - x| \leq \frac{\delta}{2}$ then

$$g(x_0) - u(x) = - \int_{\partial B_r(0)} (g(x_0) - g(y)) \frac{\partial G}{\partial n_y}(x, y) \, dy$$

$$= - \int_{\partial B_r(0) \cap |x_0 - y| \leq \delta} (g(x_0) - g(y)) \frac{\partial G}{\partial n_y}(x, y) \, dy$$

$$- \int_{\partial B_r(0) \cap |x_0 - y| > \delta} (g(x_0) - g(y)) \frac{\partial G}{\partial n_y}(x, y) \, dy$$

The first integral is easily estimated with ε
because $\int_{\partial B_r} \frac{\partial G}{\partial y_j}(x_0, y) dy = 1$.

For the second we observe that

$$\delta < |x_0 - y| \leq |x_0 - x| + |x - y| \leq \frac{\delta}{2} + |x - y|$$

so $\delta/2 < |x - y|$. It follows that the second term is bounded with

$$\begin{aligned} \dots &\leq \frac{C}{(\delta/2)^n} (r^2 - |x|^2) = \frac{C}{(\delta/2)^n} (r - |x|)(r + |x|) \leq \\ &\leq \frac{C(2r)}{(\delta/2)^n} (r - |x|) \leq \frac{C(2r)}{(\delta/2)^n} |x - x_0| \leq \varepsilon \end{aligned}$$

$$\text{if } |x - x_0| \leq \frac{\varepsilon(\delta/2)^n}{C 2r} \text{ so if}$$

$$|x - x_0| \leq \min \left\{ \frac{\delta/2}{C 2r}, \frac{\varepsilon(\delta/2)^n}{C 2r} \right\}$$

the second term is bounded by ε . \square

Problems D6, D7

↑↑

PART E: EXISTENCE FOR

$$\begin{cases} -\Delta u = f & \text{in } \Omega \\ u = g & \text{on } \partial\Omega \end{cases}$$

Motivation: if $\Delta_x G(x-y) = -\delta_{y=x}$ then

$\int_R f(y) G(r-y)$ should solve Poisson's equation because

$$\Delta_x \int_R f(y) G(x-y) = \int_R f(y) \Delta_x G(x-y) dy =$$

$$= - \int_R f(y) \delta_{y=x} dy = -f(x)$$

Therefore, we expect that $\Delta \int_R G(x-y) f(y) dy = -f(x)$.

Rigorous proof: requires careful analysis of singularity in $G(x-y)$ and some assumptions on f

As we will see, the minimal assumption will be $f \in C^2 \cap C([0,1])$.

Problem E1

There is smooth $\tilde{\chi} : \mathbb{R} \rightarrow \mathbb{R}$ s.t.

- $\tilde{\chi}(x) = 0$ for $|x| \leq 1$
- $\tilde{\chi}(x) = 1$ for $|x| \geq 2$
- $|\tilde{\chi}'(x)| \leq C \quad \forall x.$

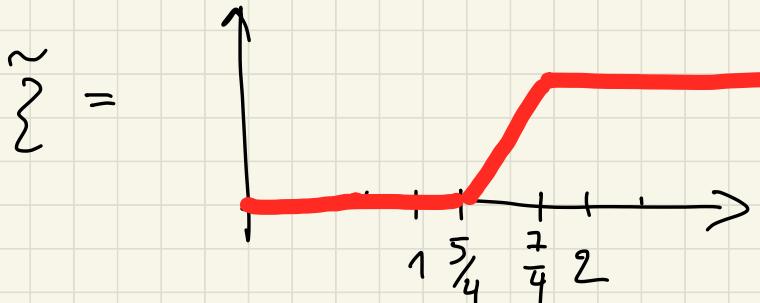
Recall mollifiers: we consider η_ε to be standard mollifying kernel i.e.

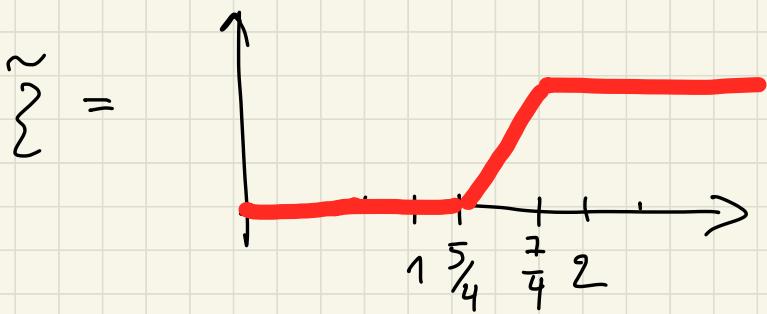
η s.t. $\int \eta = 1$, $\eta \geq 0$, $\text{supp } \eta \subset B(0, 1)$.

$$\eta_\varepsilon(x) := \frac{1}{\varepsilon^d} \eta\left(\frac{x}{\varepsilon}\right)$$

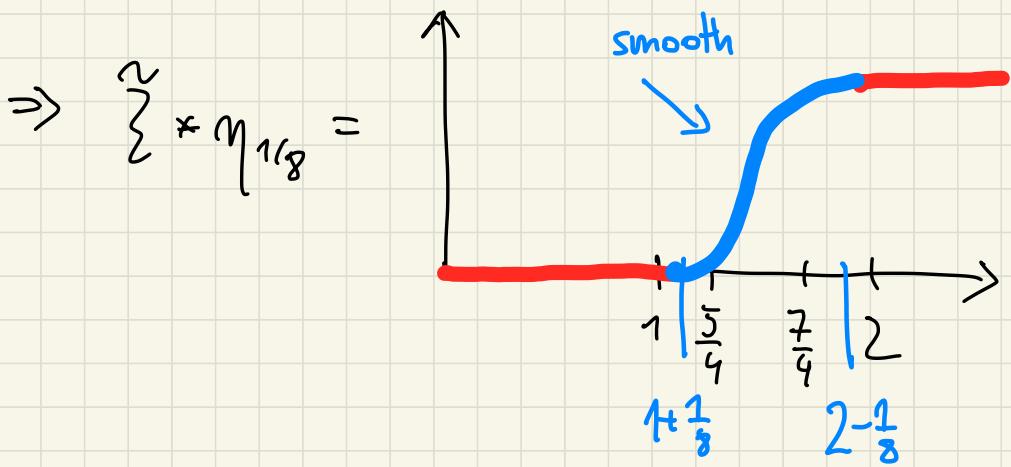
Then ($f * \eta_\varepsilon \rightarrow f$ in $C([0, 1])$, $L^p(0, 1)$ ($1 \leq p < \infty$))

We consider





Then $\tilde{\zeta} * \eta_{1/8}(x) = \int \tilde{\zeta}(y) \eta_{1/8}(x-y) dy$
 $B(x, \frac{1}{8})$



$|(\tilde{\zeta} * \eta_{1/8})'| \leq C \rightarrow$ continuous function which
 vanishes outside $[1, 2]$.

Problem E2 $\sum_{\varepsilon}(x) = \sum\left(\frac{|x|^2}{\varepsilon^2}\right)$

Then $\sum_{\varepsilon}(x) = 0$ when $|x| \leq \varepsilon$

$\sum_{\varepsilon}(x) = 1$ when $|x| \geq \sqrt{2} \varepsilon$

$$D \sum_{\varepsilon}(x) = \sum\left(\frac{|x|^2}{\varepsilon^2}\right) \cdot \frac{1}{\varepsilon^2} 2x$$

$$= \frac{2x}{\varepsilon} \cdot \sum\left(\frac{|x|^2}{\varepsilon^2}\right) \cdot \frac{1}{\varepsilon} \leq \frac{C}{\varepsilon}$$

$\leq \sqrt{2}.$

Problem E3

Claim: $w_f(x) := \int_{\Omega} \phi(x-y) f(y) dy$

$$w_f^\varepsilon(x) := \int_{\Omega} \phi(x-y) \sum_{\varepsilon} (x-y) f(y) dy$$

$w_f^\varepsilon \rightarrow w_f$ uniformly in Ω

$$\text{Proof: } |w_f(x) - w_f^\varepsilon(x)| = \left| \int_{\mathbb{R}} \Phi(x-y) f(y) \underbrace{\left[1 - \zeta^\varepsilon(x-y) \right]}_{=0 \text{ when }} dy \right|$$

$$\zeta^\varepsilon(x-y) = 1 \text{ when}$$

$$|x-y| \geq \sqrt{2}\varepsilon$$

$$\leq \left| \int_{|x-y| \leq \sqrt{2}\varepsilon} \underbrace{\Phi(x-y)}_{\frac{C}{|x-y|^{n-2}}} f(y) dy \right| \leq \left| \int_{|x-y| \leq \sqrt{2}\varepsilon} \frac{C}{|x-y|^{n-2}} f(y) dy \right|$$

$$\leq \int_{B(x, \sqrt{2}\varepsilon)} \frac{C}{|x-y|^{n-2}} |f(y)| dy \leq C \|f\|_\infty \int_0^{\sqrt{2}\varepsilon} \frac{1}{r^{n-2}} r^{n-1} dr$$

$$\leq C \|f\|_\infty \varepsilon^2 \rightarrow 0$$

Case $d=2 \Rightarrow$ SIMILAR $\left(\int_0^\varepsilon r \log r dr \right)$

Problem E4

Target: compute derivative of $\int_{\mathbb{R}} \hat{\phi}(x-y) f(y) dy$

$$\text{Claim: } D^i w_f(x) = \int_{\mathbb{R}} D^i \hat{\phi}(x-y) f(y) dy$$

Step 1: this function is well-defined, i.e. for all $x \in \mathbb{R}$ $D^i w_f(x) < \infty$. Proof: find R_x s.t. $\mathbb{R} \subset B_{R_x}(x)$ so that

$$\begin{aligned} |D^i w_f(x)| &\leq \int_{B_{R_x}(x)} |D^i \hat{\phi}(x-y)| \|f\|_\infty dy \\ &\leq \|f\|_\infty \int_0^{R_x} \frac{C}{r^{n-1}} r^{n-1} dr \leq (\|f\|_\infty R_x) < \infty \end{aligned}$$

$$\begin{aligned} \underline{\text{Step 2: }} \quad D^i w_f^\varepsilon(x) &= \int_{\mathbb{R}} D^i [\hat{\phi}(x-y) \sum \varepsilon(x-y)] f(y) dy \\ &= \int_{\mathbb{R}} D^i \hat{\phi}(x-y) f(y) dy + \int_{\mathbb{R}} \hat{\phi}(x-y) D^i [\sum \varepsilon(x-y)] f(y) dy \end{aligned}$$

how to differentiate integrals:

$$\frac{d}{dx} \int_{\Omega} F(x,y) dy \stackrel{?}{=} \int_{\Omega} \frac{d}{dx} F(x,y) dy$$

? (*)

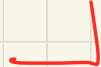
$$\frac{1}{h} \left[\int_{\Omega} F(x+h,y) dy - \int_{\Omega} F(x,y) dy \right] =$$

$$= \int_{\Omega} \frac{F(x+h,y) - F(x,y)}{h} dy$$

\downarrow
 $\frac{d}{dx} F(x,y)$ pointwise

Note $\frac{F(x+h,y) - F(x,y)}{h} \leq \|F_x\|_{\infty}$ so if

We integrate over the set of finite measure and derivative is bounded we are allowed to conclude (*).



Second integral :

$$\left| \int_{\mathbb{R}} D^i \phi(x-y) D^j \{^\varepsilon(x-y) f(y) dy \} \right| \leq$$

$$\leq \frac{1}{\varepsilon} \int |D^i \phi(x-y)| |f(y)| dy \leq \|f\|_\infty \frac{1}{\varepsilon} \int \frac{C}{|x-y|^{n-2}}$$

$\varepsilon \leq |x-y| \leq \sqrt{2} \varepsilon$

$$\leq C \|f\|_\infty \frac{1}{\varepsilon} \int_{\varepsilon}^{\sqrt{2}\varepsilon} \frac{r^{n-1}}{r^{n-2}} dr = C \|f\|_\infty \varepsilon \rightarrow 0$$

uniformly in x

First integral:

$$\left| \int_{\mathbb{R}} D^i \phi(x-y) [\{^\varepsilon(x-y) - 1] f(y) dy \right| \leq$$

$$\leq \int_{|x-y| \leq \sqrt{2}\varepsilon} |D^i \phi(x-y)| \|f\|_\infty dy \leq$$

$$\leq C \|f\|_\infty \int_{|x-y| \leq \sqrt{2}\varepsilon} \frac{1}{|x-y|^{n-1}} dy \leq C \|f\|_\infty \int_0^{\sqrt{2}\varepsilon} \frac{1}{r^{n-1}} r^{n-1} dr$$

$$= C \|f\|_\infty \sqrt{2} \varepsilon \rightarrow 0$$

again, uniformly in $x \in \Omega$.

Conclusion:

$$w_f^\varepsilon \xrightarrow{\quad} w_f \quad \text{in } \Omega$$

$$D_i w_f^\varepsilon \xrightarrow{\quad} \int_{\Omega} D_i \Phi(x-y) f(y) dy \quad \text{in } \Omega$$

\Rightarrow for all $\tilde{\Omega}$ compact, $\tilde{\Omega} \subset \Omega$, w_f^ε is Cauchy in $C^1(\tilde{\Omega})$ = Banach space.

$$\Rightarrow w_f^\varepsilon \rightarrow g \quad \text{in } C^1(\tilde{\Omega}) \Rightarrow g = w_f \Rightarrow$$

$$w_f \in C^1(\tilde{\Omega})$$

$$D_i w_f^\varepsilon \rightarrow D_i w_f \quad \text{in } C^1(\tilde{\Omega}) \Rightarrow$$

$$D_i w_f = \int_{\Omega} D_i \Phi(x-y) f(y) dy$$

Problem E5

How to compute 2nd derivative?

$$D_{ij} w_f(x) = \int_{\mathbb{R}} D_{ij} \hat{\phi}(x-y) f(y) dy$$

?

This is not well-defined for $f \in L^\infty$.

$$|D_{ij} \hat{\phi}(x-y)| \leq \frac{C}{|x-y|^n}$$

$$\text{But } \int_{\mathbb{R}} \frac{C}{|x-y|^n} dy \leq \int_{B_R(x)} \frac{C}{|x-y|^n} dy \approx \int_0^R \frac{1}{r^n} r^{n-1} dr = \infty$$

This motivated looking for conditions on f allowing to compute $D_{ij} w_f$. The condition is Hölder continuity

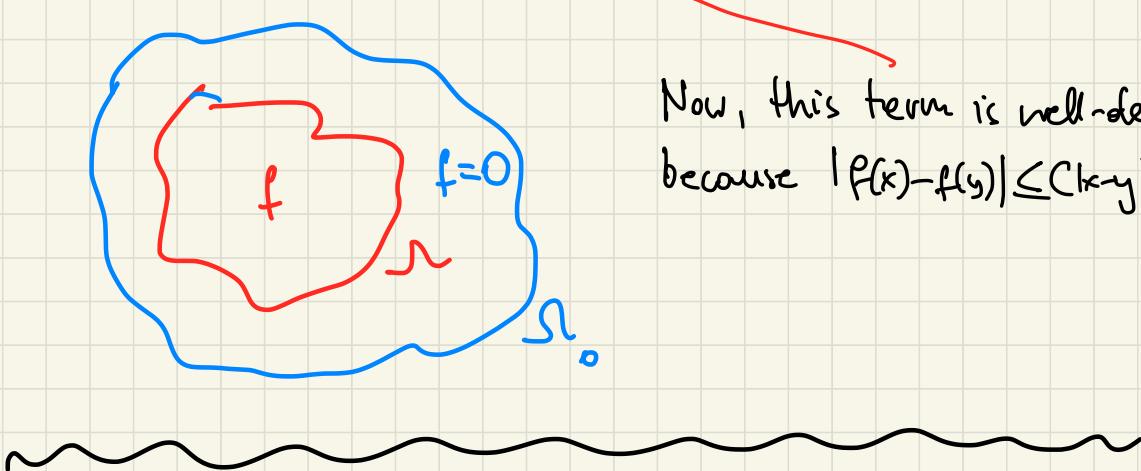
i.e. $\exists C \quad |f(x) - f(y)| \leq C|x-y|^\alpha$.

↑:

E4. Let $f \in L^\infty(\Omega) \cap C^\alpha(\Omega)$ for some $\alpha \in (0, 1]$. Consider Ω_0 such that $\Omega \subset \Omega_0$ and extend $f = 0$ in $\Omega_0 \setminus \Omega$. Then $w_f \in C^2(\Omega)$ and

$$D_{i,j} w_f(x) = \int_{\Omega_0} D_{i,j} \Phi(x-y) (f(x) - f(y)) dy - f(x) \int_{\partial\Omega_0} D_i \Phi(x-y) n_j(y) dS(y)$$

where n_j is the j -th component of \mathbf{n} .



Schauder estimates for $-\Delta$: $(-\Delta u = f)$

$$f \in C^0 \quad \cancel{\Rightarrow} \quad u \in C^2$$

$$f \in C^1 \quad \Rightarrow \quad u \in C^{2+\perp} \quad (\text{Schauder})$$

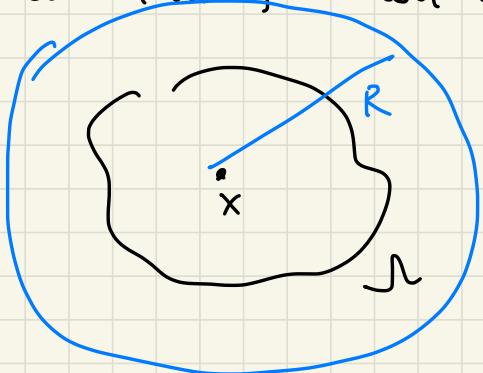
$$f \in L^2 \quad \Rightarrow \quad u \in H^2 \quad (\text{future})$$

(0 weak derivatives) (two weak derivatives)

Problem E6

$$D_{i,j} u_f(x) = \int_{\Omega_0} D_{i,j} \Phi(x-y) (f(x) - f(y)) dy \\ - f(x) \int_{\partial \Omega_0} D_i \Phi(x-y) u_j(y) dS(y)$$

and this function is well-defined.



$$\Omega \subset B_R(x) := \Omega_0$$

$$D_i \Phi(x) = \frac{x_i}{n \det|x|^n}.$$

$$\sum_{i=1}^n D_{i,i} u_f(x) = - \sum_{i=1}^n \int_{\partial B_R(x)} D_i \Phi(x-y) u_i(y) dS(y) \\ = - f(x) \iint_{\partial B_R(x)} \frac{x_i - y_i}{n \det|x-y|^n} u_i(y) dS(y) = \\ = - f(x) \int_{\partial B_R(0)} \frac{\langle y_i, n \rangle}{n \det|y|^n} dS(y)$$

$$\begin{aligned}
 & -f(x) \int_{\partial B_R(0)} \frac{\langle y_1^n \rangle}{n d_m R^n} dS(y) = \\
 &= -f(x) \frac{1}{n d_m R^n} \int_{\partial B_R(0)} \langle y_1^n \rangle dS(y) = \\
 &= -f(x) \frac{1}{n d_m R^n} \int_{B_n(0)} \operatorname{div}(y) dx = -f(x).
 \end{aligned}$$

□.

Problem E7 In the lecture $\Omega = \mathbb{R}^n$ and $f \in C_c^\infty$

then $\int_{\mathbb{R}^n} \Phi(x-y) f(y) dy$ solves Poisson equation,
 $\|$

$$\int_{\mathbb{R}^n} \Phi(y) f(x-y) dy$$

$f \in C_c^\infty$ so \uparrow can be directly differentiated.

Note that if $f \in C^\infty$ and f is comp. supp. we can write

$$\int_{\mathbb{R}^n} \Phi(y) f(x-y) dy = \int_{\text{supp } f} \Phi(x-y) f(y) dy$$

and we can use our theory to see that $u \in C^2$ (even $C^{2+\alpha}$ by Schauder estimate).

Problem E8 Uniqueness to $\begin{cases} -\Delta u = f \\ u = g \end{cases}$ is

clear. For existence solve

$$\begin{cases} -\Delta v = 0 & \text{in } \Omega \\ v = g - w_f & \text{on } \partial\Omega \end{cases} \text{ then } u = v + w_f \text{ solves.}$$